

# Simulation of Fourth-Order Laterally-Coupled Gratings

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# Outline

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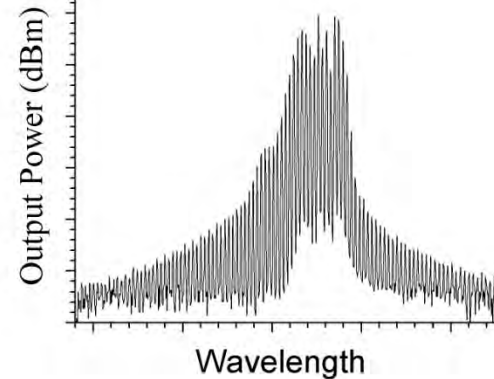
- Laterally-coupled distributed feedback (LC-DFB) laser introduction
- Modified coupled-mode theory
- Design parameters analysis (Comsol simulations)
- Conclusions



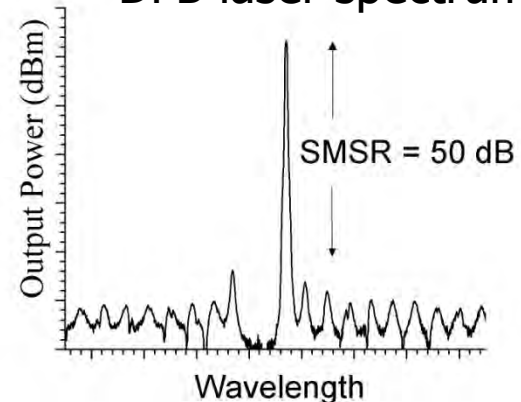
# DFB laser introduction

- Distributed feedback (DFB) lasers use a wavelength-selective grating for improved spectral purity
- Lower temperature and wavelength sensitivity than Fabry-Perot designs
- Better longitudinal mode selectivity; reduces multiple cavity modes of Fabry-Perot lasers
- Used for telecommunications, particularly DWDM, where there is a requirement for low linewidth and good stability

FP laser spectrum



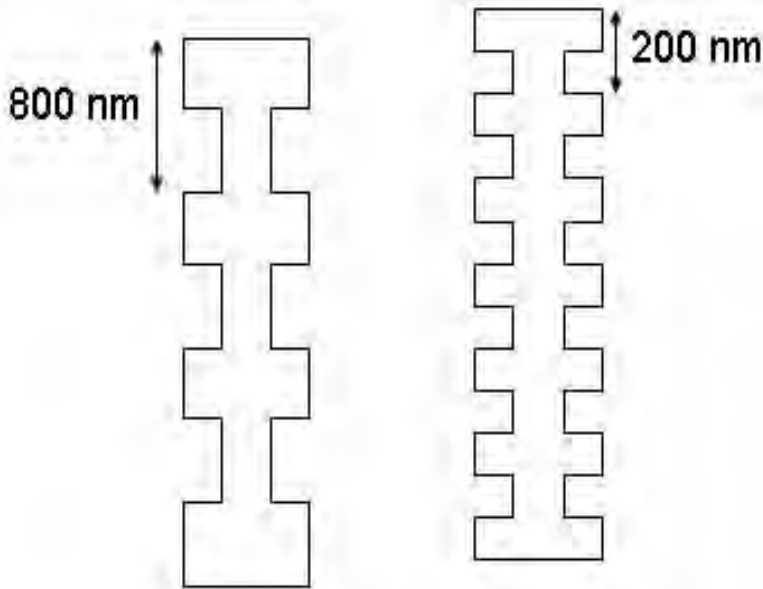
DFB laser spectrum





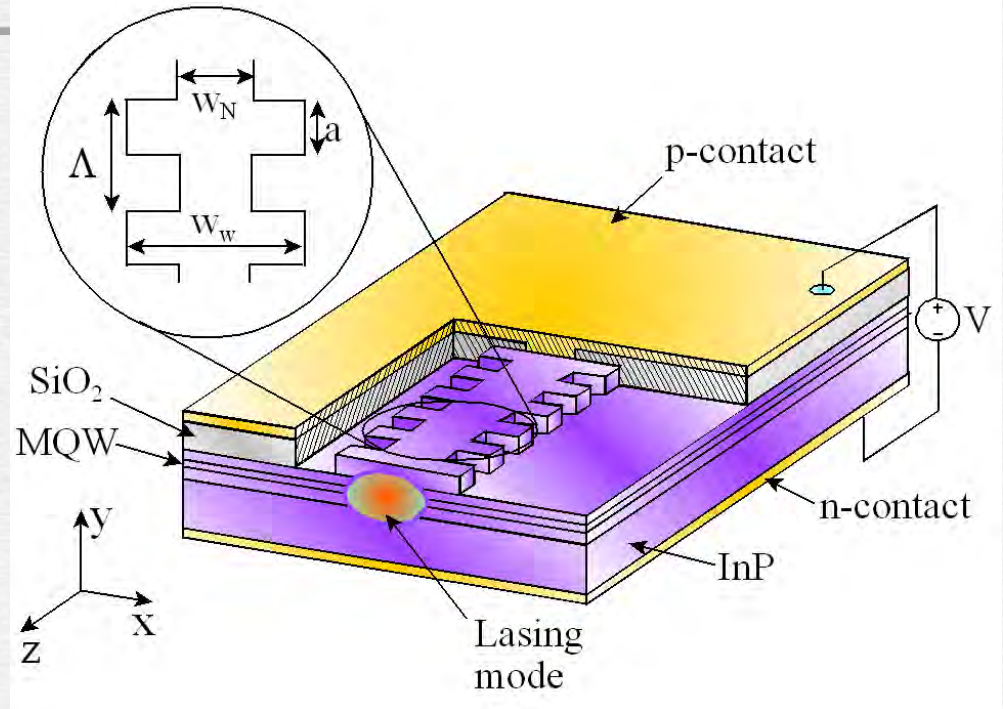
# LC-DFB laser introduction

- Grating patterned out of upper ridge waveguide: No re-growth is needed
- Can be fabricated using stepper lithography or nano-imprinting – amenable to mass-manufacturing



Fourth-order

First-order



- Some Higher-order gratings advantages:
  - ease up the stringent fabrication limitations
  - allow better longitudinal mode discrimination



# Modified coupled-mode theory

In higher order gratings, additional terms are included to account for light radiating in transverse direction:

$$\begin{aligned}\frac{dA}{dz} + (-\alpha - i\delta - i\zeta_1)A &= i\left(\kappa_p^* + \zeta_2\right)B \\ -\frac{dB}{dz} + (-\alpha - i\delta - i\zeta_3)B &= i\left(\kappa_p + \zeta_4\right)A\end{aligned}$$

A,B = longitudinal mode fields

$\kappa_p$  = Coupling coefficient

$\alpha$  = modal gain

$\delta$  = Bragg frequency detuning

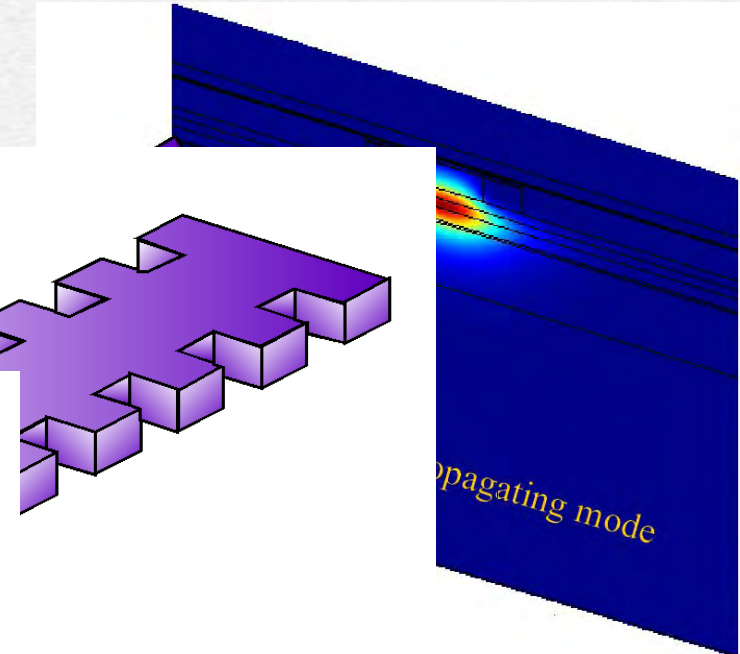
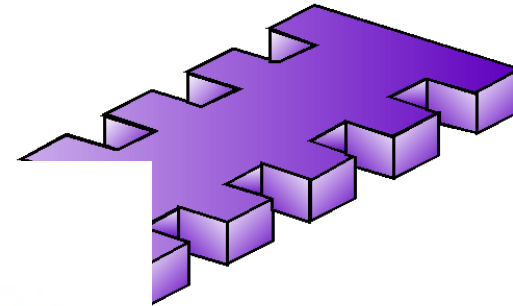
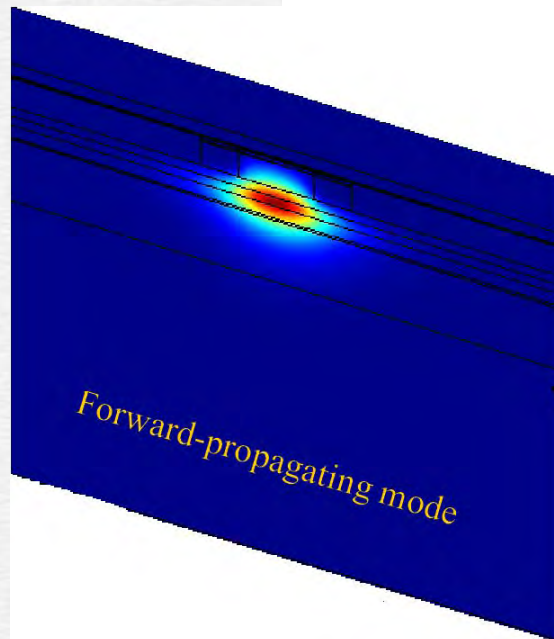
$\zeta_{1,\dots,4}$  = Streifer correction terms



# Coupling coefficient

$$\frac{dA}{dz} + (-\alpha - i\delta - i\zeta_1)A = i\left(\kappa_p^* + \zeta_2\right)B$$

$$-\frac{dB}{dz} + (-\alpha - i\delta - i\zeta_3)B = i\left(\kappa_p + \zeta_4\right)A$$



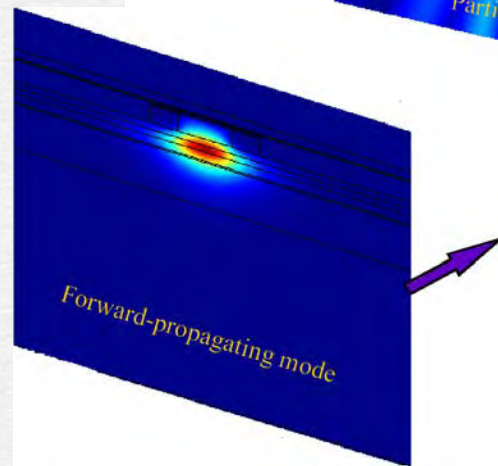
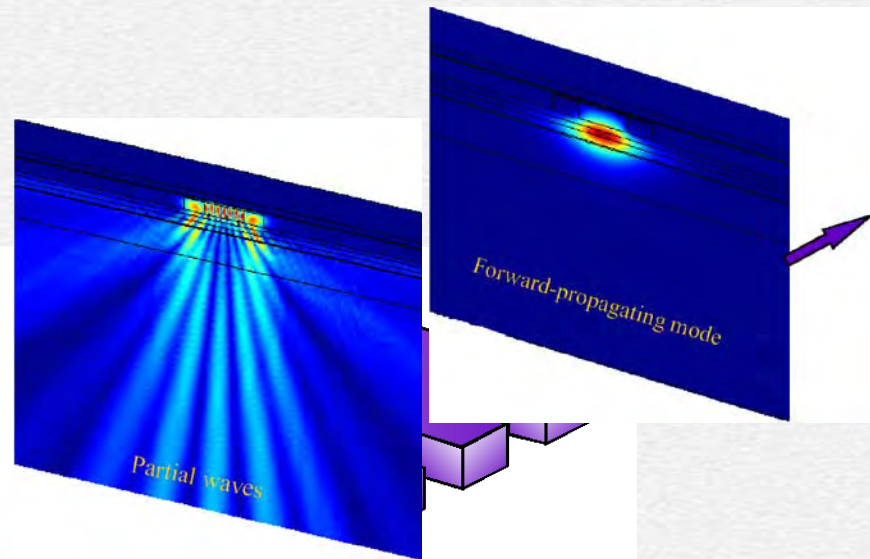
Coupling coefficient,  $\kappa_p$ , measures amount of coupling between forward- and backward-propagating fundamental modes



# Streifer correction terms

$$\frac{dA}{dz} + \left( -\alpha - i\delta - i\zeta_1 \right) A = i \left( \kappa_p^* + \zeta_2 \right) B$$

$$-\frac{dB}{dz} + \left( -\alpha - i\delta - i\zeta_3 \right) B = i \left( \kappa_p \right) A$$



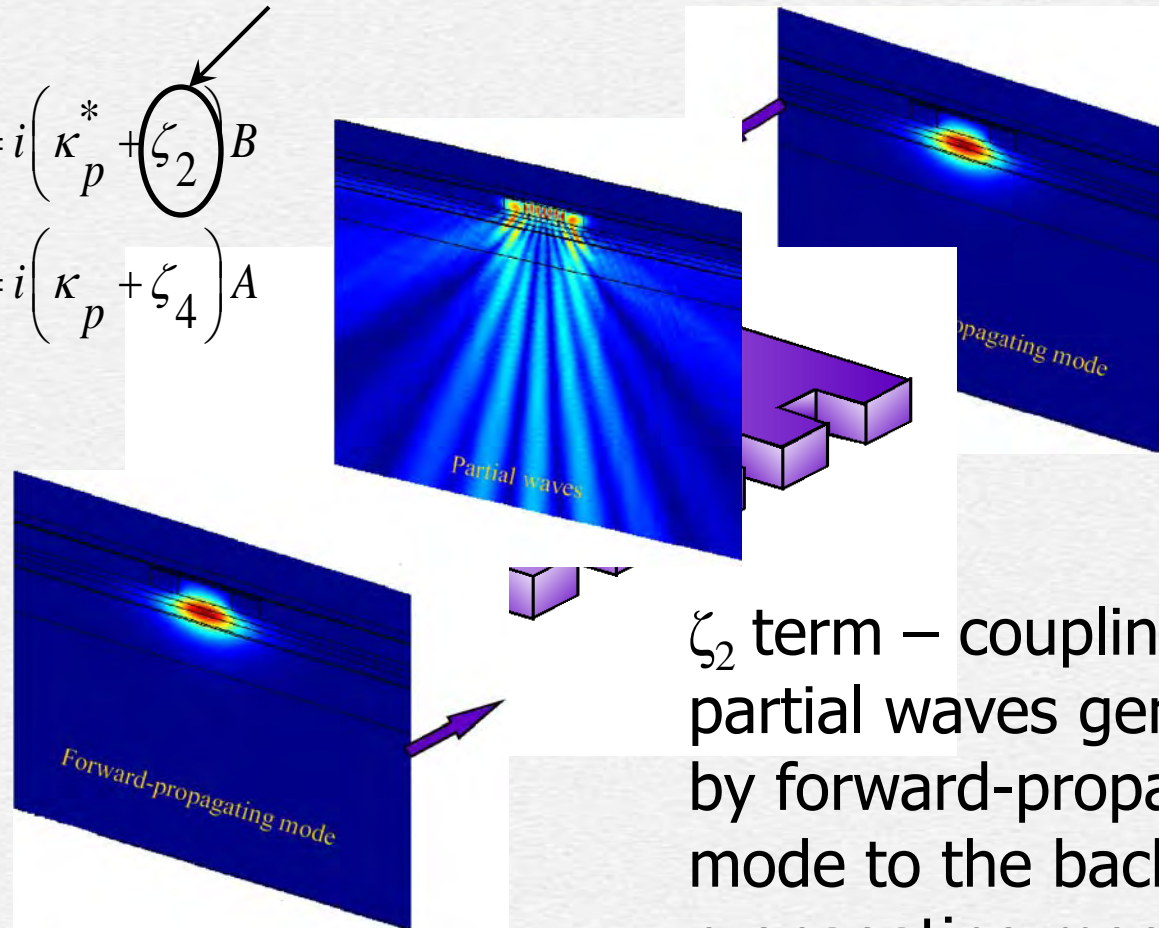
$\zeta_1$  term – coupling of partial waves generated by forward-propagating mode to the forward-propagating mode



# Streifer correction terms

$$\frac{dA}{dz} + (-\alpha - i\delta - i\zeta_1)A = i\left(\kappa_p^* + \zeta_2\right)B$$

$$-\frac{dB}{dz} + (-\alpha - i\delta - i\zeta_3)B = i\left(\kappa_p + \zeta_4\right)A$$



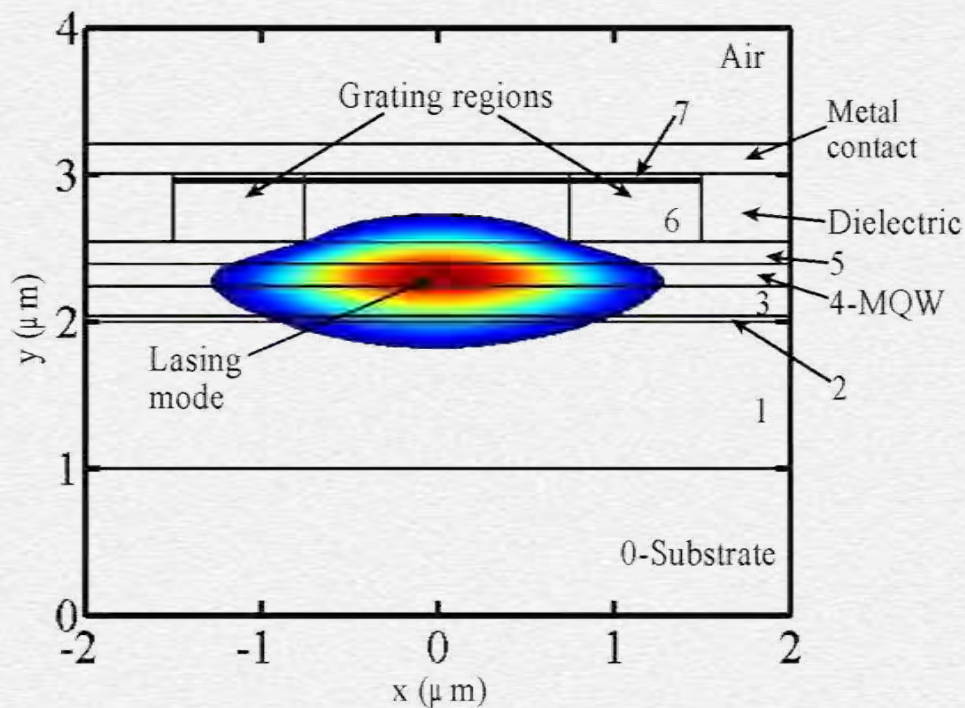
$\zeta_2$  term – coupling of partial waves generated by forward-propagating mode to the backward-propagating mode





# Modified coupled-mode theory

- Quasi-TE fundamental mode  $\epsilon_0(x,y)$  is evanescently coupled to laterally-positioned grating region
- MQW active region
- Au/Pt/Ti contact with  $\text{SiO}_2$  dielectric





# Modified coupled-mode theory

Correction terms are determined through the solution of the driven wave equation having as a source the quasi-TE fundamental mode field  $\varepsilon_0(x, y)$ :

$$\frac{\partial^2 \varepsilon_m^{(i)}(x, y)}{\partial x^2} + \frac{\partial^2 \varepsilon_m^{(i)}(x, y)}{\partial y^2} + [k_0^2 n_0^2(x, y) - \beta_m^2] \varepsilon_m^{(i)}(x, y) = -k_0^2 A_{m-i}(x, y) \varepsilon_0(x, y), \quad m \neq i, i = 0, p.$$

$\varepsilon_m$  = partial wave field of order  $m$

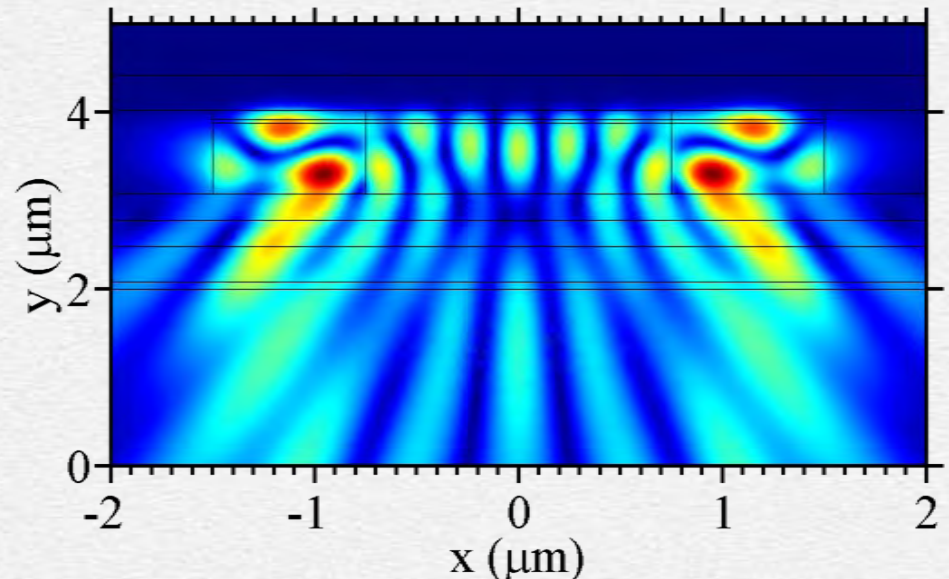
$\varepsilon_0$  = fundamental TE mode field

$k_0$  = Vacuum wavenumber

$\beta_m$  = partial wave propagation constant

$A_q$  =  $q^{\text{th}}$  order Fourier coefficient

Radiating partial wave fields are calculated using the FEM with absorbing boundary conditions

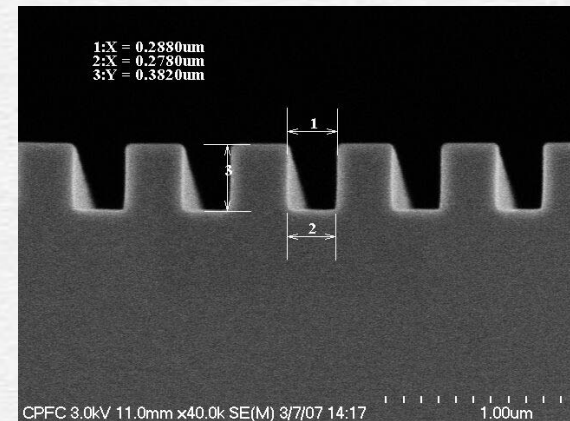
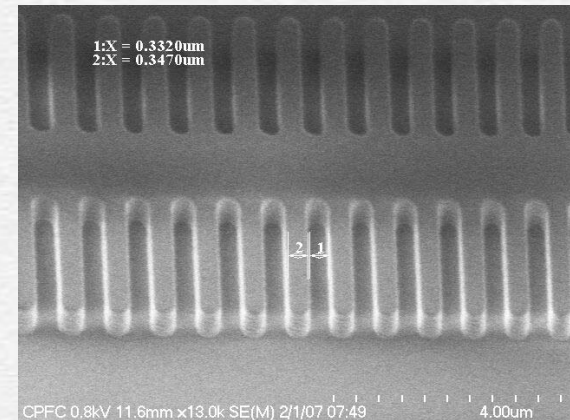


# Grating strength

- A measure of grating strength in higher order gratings is the effective coupling coefficient:

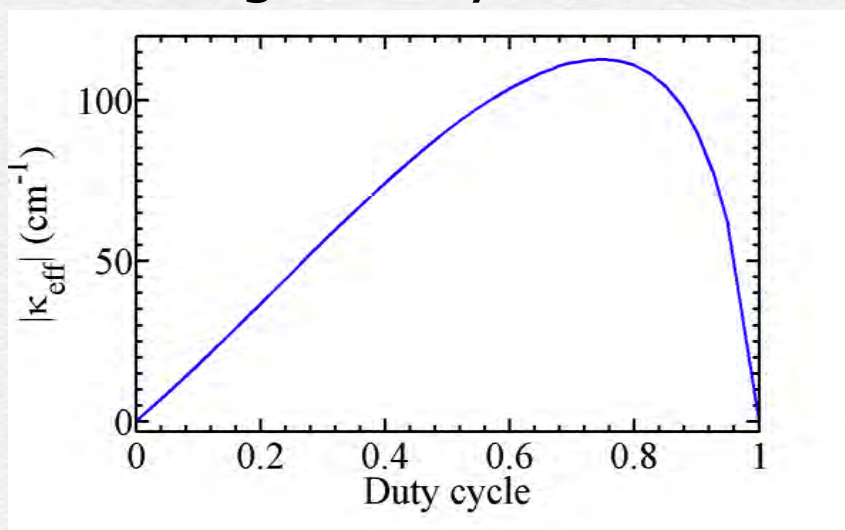
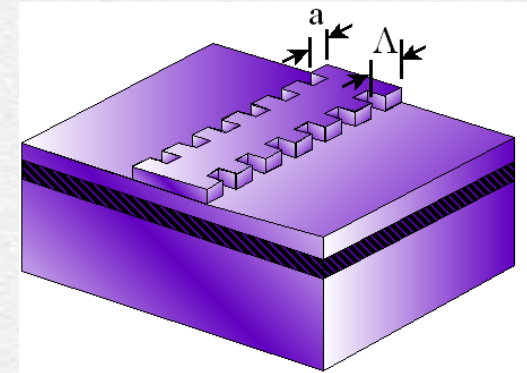
$$\kappa_{eff} = \sqrt{(\kappa_p^* + \zeta_2)(\kappa_p + \zeta_4)} = |\kappa_{eff}| e^{j\phi(\kappa_{eff})}$$

- Combination of all coupling terms between forward- to backward- propagating (and vice versa) waves
- Values of  $\kappa L \approx 1.25$  (L=cavity length) are desirable for DFB lasers

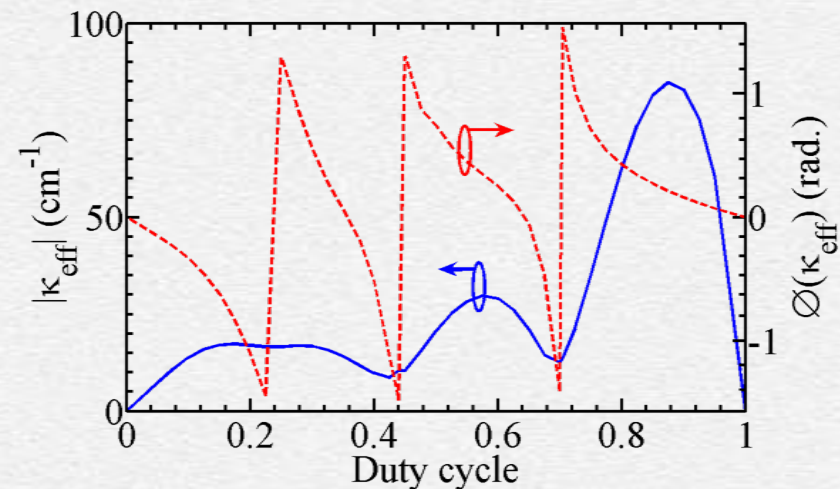


# Duty cycle/Grating order

- Duty cycle =  $a/\Lambda$
- LC-DFB performance is sensitive to duty cycle: larger values ( $> 0.5$ ) are generally better



First-order



Fourth-order

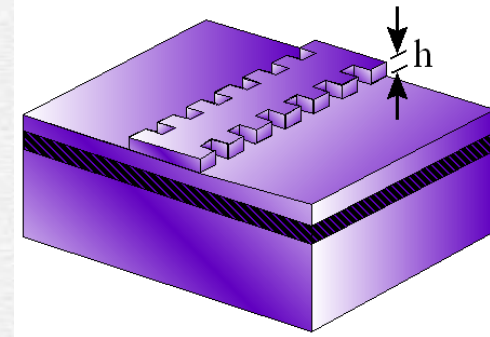


# Grating height: fixed duty cycle of 0.6

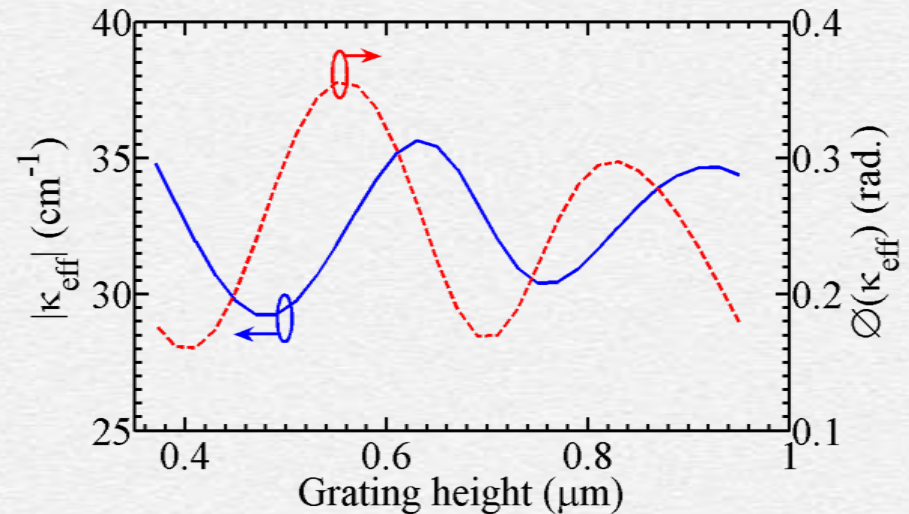
- Change in grating height produces a periodic change in effective coupling coefficient due to correction terms – resonance effects between upper metal contact and lower dielectric layers

- Uncorrected  $\kappa_p$  at a grating height of  $0.47 \mu\text{m}$  is  $23.3 \text{ cm}^{-1}$  ( $|\kappa_{\text{eff}}| \sim 29 \text{ cm}^{-1}$ ); this underscores the importance of including partial wave terms in the calculations

- First peak in  $|\kappa_{\text{eff}}|$  at a height of  $\sim 0.64 \mu\text{m}$  represents a good stable fabrication point

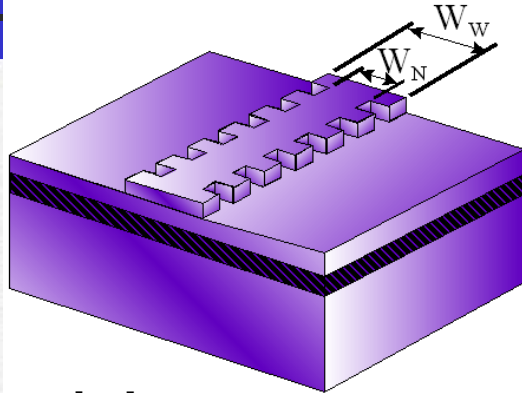


$$W_w/W_n = 3\mu\text{m}/1.5\mu\text{m}$$

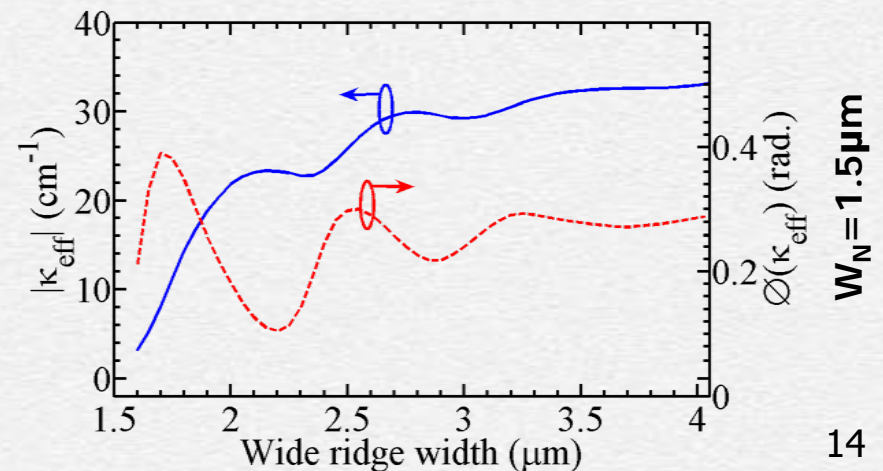
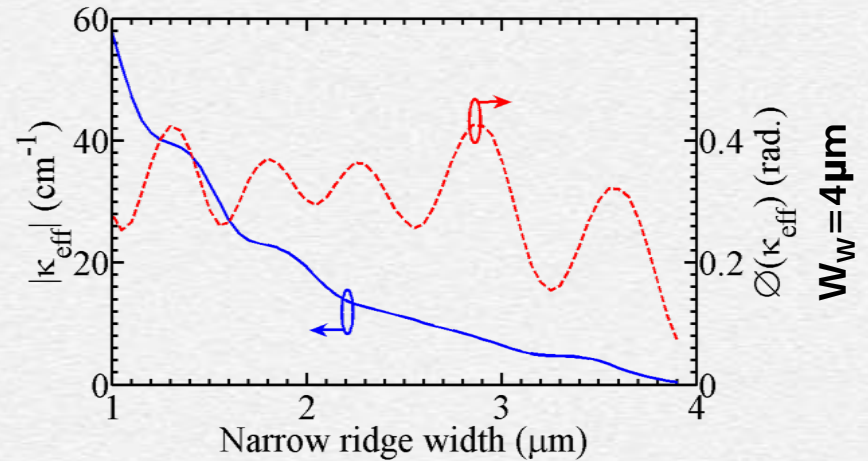




# Grating width: fixed duty cycle of 0.6



- More subtle resonances are observed with grating width
- Less reflection from sidewalls of grating than at metal interface
- Coupling increases for a more pronounced grating (i.e. width ratio  $(W_W/W_N) \gg 1$ )
- Good target fabrication window:  
 $W_W > 3 \mu\text{m}$   $W_N < 1.5 \mu\text{m}$





# Conclusions

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- LC-DFB lasers with higher order gratings can be effectively modeled using COMSOL Multiphysics
- Laser performance and tolerances are determined by grating geometry, including duty cycle, grating width and grating height
- Radiating partial wave effects should be included in the calculation of LC-DFB lasers with higher order gratings



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